



Third International Conference on Quantum Error Correction • 15-19 December 2014 • Zurich, Switzerland

# A General Transfer-Function Approach to Noise Filtering in Open-Loop Quantum Control

**Lorenza Viola**

**Dept. Physics & Astronomy  
Dartmouth College**

Paz-Silva & LV, arXiv:1408.3836, Phys. Rev. Lett. (2014) [in press]





**Third International Conference on Quantum Error Correction • 15-19 December 2014 • Zurich, Switzerland**



**Gerardo Paz-Silva**  
Dartmouth



**Michael Biercuk  
& Todd Green**  
U. Sydney



**Ken Brown &  
Chingiz Kabytaev**  
GeorgiaTech



**Will Oliver**  
MIT



# Motivation

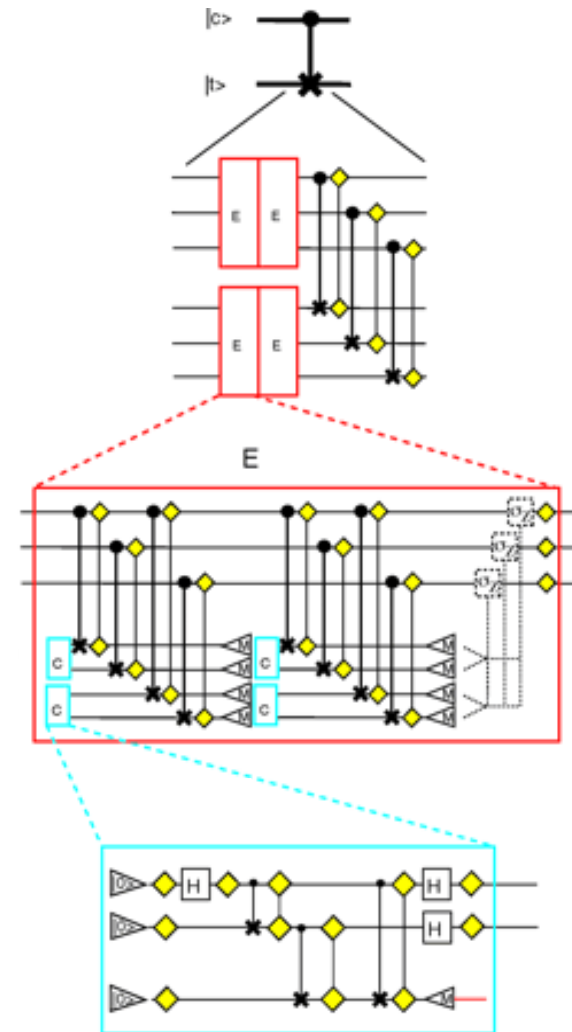
Goal: High-precision, robust control of *realistic* quantum-dynamical systems.

- Real-world quantum control systems typically entail:

- *Noisy*, irreversible open-system dynamics...
- *Imperfectly characterized* dynamical models...
- *Limited* control resources...
- ⋮

- **Broad significance across coherent quantum sciences:**

- High-resolution imaging and spectroscopy...
- Quantum chemistry and biology...
- Quantum metrology, sensing and identification...
- **High-fidelity QIP, fault-tolerant QEC...**
- ⋮



# Motivation

Goal: High-precision, robust control of *realistic* quantum-dynamical systems.

- Real-world quantum control systems typically entail:

- *Noisy*, irreversible open-system dynamics...
- *Imperfectly characterized* dynamical models...
- *Limited* control resources...

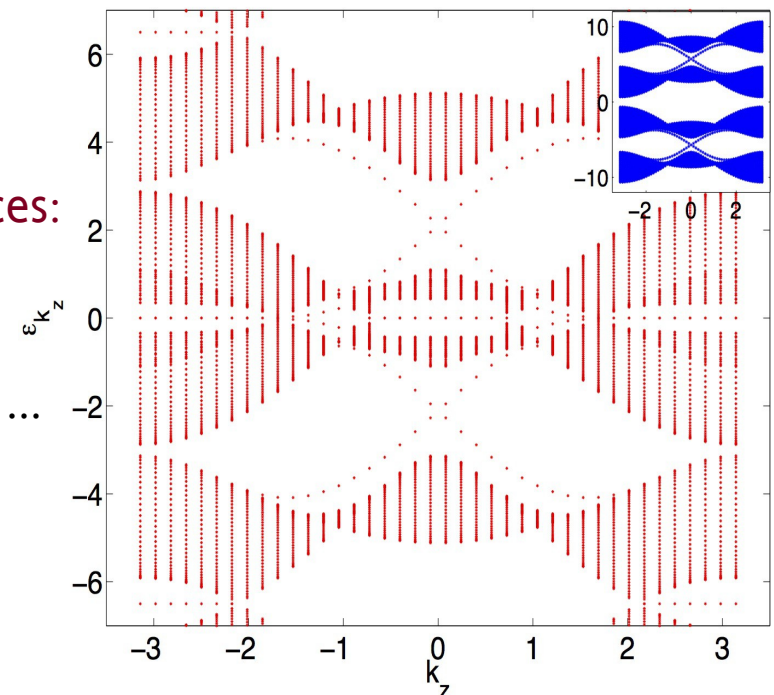
⋮

- **Broad significance across coherent quantum sciences:**

- High-resolution imaging and spectroscopy...
- Quantum chemistry and biology...
- Quantum metrology, sensing and identification ...
- High-fidelity QIP, fault-tolerant QEC...
- Engineering of novel quantum matter...

⋮

Poudel, Ortiz & LV,  
Floquet Majorana flat bands,  
ArXiv:1412.2639



# The premise: Dynamical QEC

Open-loop *Hamiltonian engineering* [both *closed* and *open* systems]:  
Dynamical control solely based on *unitary* control resources.

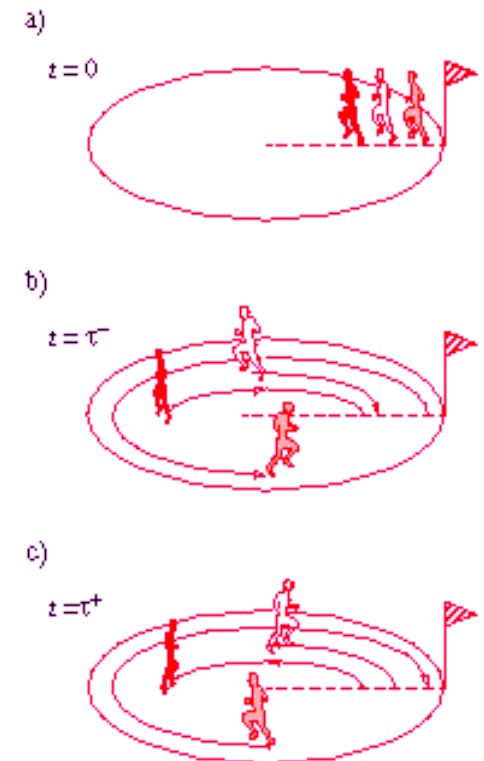
Simplest setting: Multi-pulse decoherence control for quantum memory  $\Rightarrow$  DD

LV & Lloyd, PRA 1998.

Key principle: Time-scale separation  $\Rightarrow$  'Coherent averaging'

Paradigmatic example: Spin echo  $\Leftrightarrow$  Effective time-reversal

Hahn, PR 1950.



# The premise: Dynamical QEC

Open-loop *Hamiltonian engineering* [both *closed* and *open* systems]:  
Dynamical control solely based on *unitary* control resources.

Simplest setting: Multi-pulse decoherence control for quantum memory  $\Rightarrow$  DD

LV & Lloyd, PRA 1998.

Key principle: Time-scale separation  $\Rightarrow$  'Coherent averaging'

Paradigmatic example: Spin echo  $\Leftrightarrow$  Effective time-reversal

Hahn, PR 1950.

Key features: 'Non-Markovian' quantum dynamics

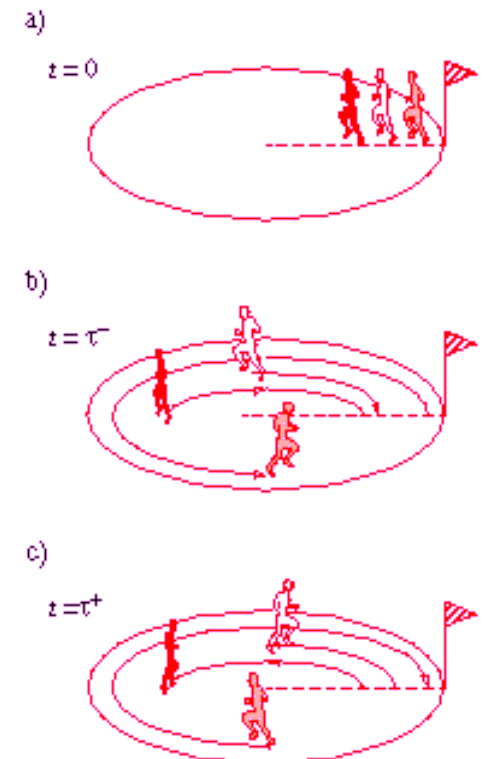
(1) Dynamical error suppression is achieved in a *perturbative* sense

$$\frac{\tau_{\text{ctrl}}}{\tau_{\text{corr}}} \sim \omega_c \tau_{\text{ctrl}} \quad \text{small parameter}$$

(2) Unwanted dynamics may include coupling to *quantum* bath

(3) Dynamical QEC is achievable *without* requiring full/quantitative knowledge of error sources

[ $\Rightarrow$  built-in robustness against 'model uncertainty']



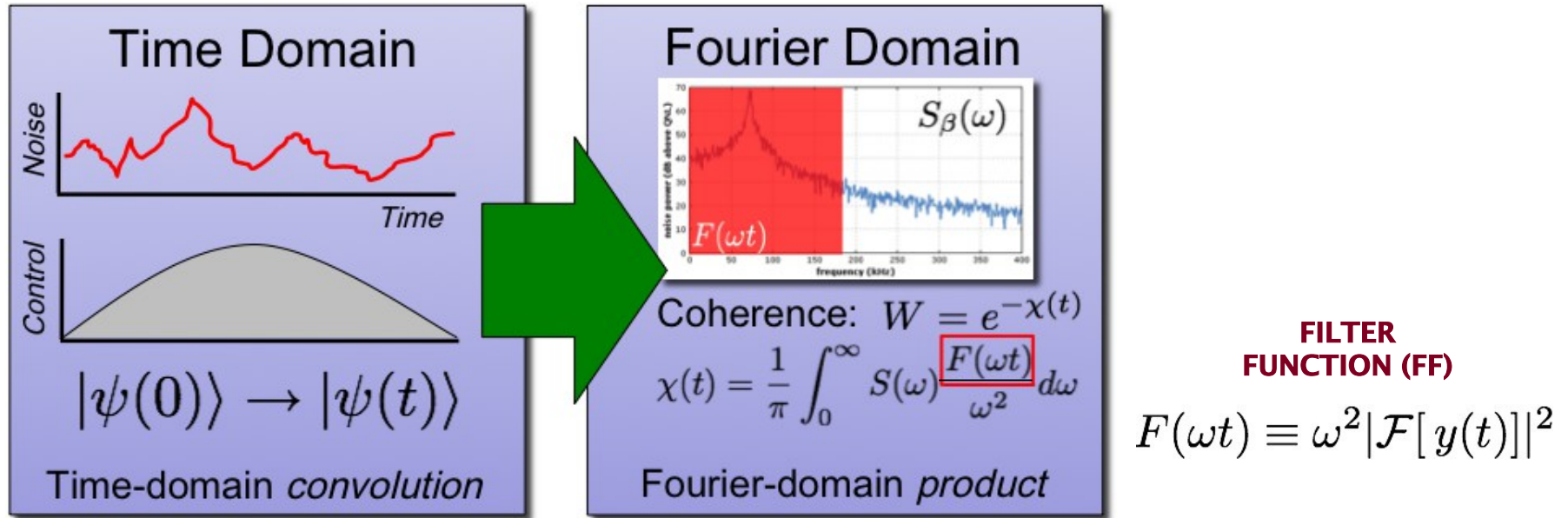
# Quantum control tasks

---

- Hamiltonian engineering techniques provide a versatile tool for *dynamical control* and *physical-layer* decoherence suppression in a variety of QIP settings:
  - Arbitrary state preservation  $\Rightarrow$  DQEC for *quantum memory*
    - ✓ Pulsed DD – 'Bang-Bang' (BB) limit/instantaneous pulses
    - ✓ Pulsed DD – Bounded control ('Eulerian')/'fat' pulses
    - ✓ Continuous-(Wave, CW) [always-on] DD
  - Quantum gate synthesis  $\Rightarrow$  DQEC for *quantum computation*
    - ✓ Hybrid DD-QC schemes – BB, w or w/o encoding
    - ✓ Dynamically corrected gates (DCGs) – Bounded control only
    - ✓ Composite pulses – Bounded control only
  - Quantum system identification  $\Rightarrow$  Dynamical control for *signal/noise estimation*
    - ✓ Signal reconstruction – dynamic parameter estimation ('Walsh spectroscopy')
    - ✓ Spectral reconstruction – DD noise spectroscopy
  - Hamiltonian synthesis  $\Rightarrow$  Dynamical control for *quantum simulation*
    - ✓ Closed-system [many-body, BB and Eulerian] Hamiltonian simulation
    - ✓ Open-system [dynamically corrected] Hamiltonian simulation
    - ⋮

# Time vs frequency domain: Filter transfer functions

Kurizki *et al*/PRL 2001; Uhrig PRL 2007; Cywinski *et al*, PRB 2008; Khodjasteh *et al*, PRA 2011; Biercuk *et al*, JPB 2011; Hayes *et al*, PRA 2011; Green *et al*, PRL 2012, NJP 2013; Kabytayev *et al*, PRA 2014...



- Picture the control modulation as enacting a 'noise filter' in frequency domain:

→ Simplest case: Single qubit under *classical Gaussian dephasing*, DD via perfect  $\pi$  pulses

$$U_{\text{ctrl}}(t) = \sigma_x^{(y(t)+1)/2} = \begin{cases} \mathbb{I} \\ \sigma_x \end{cases} \quad F(\omega t) \propto (\omega t)^{2(\delta+1)}, \quad \omega \rightarrow 0$$

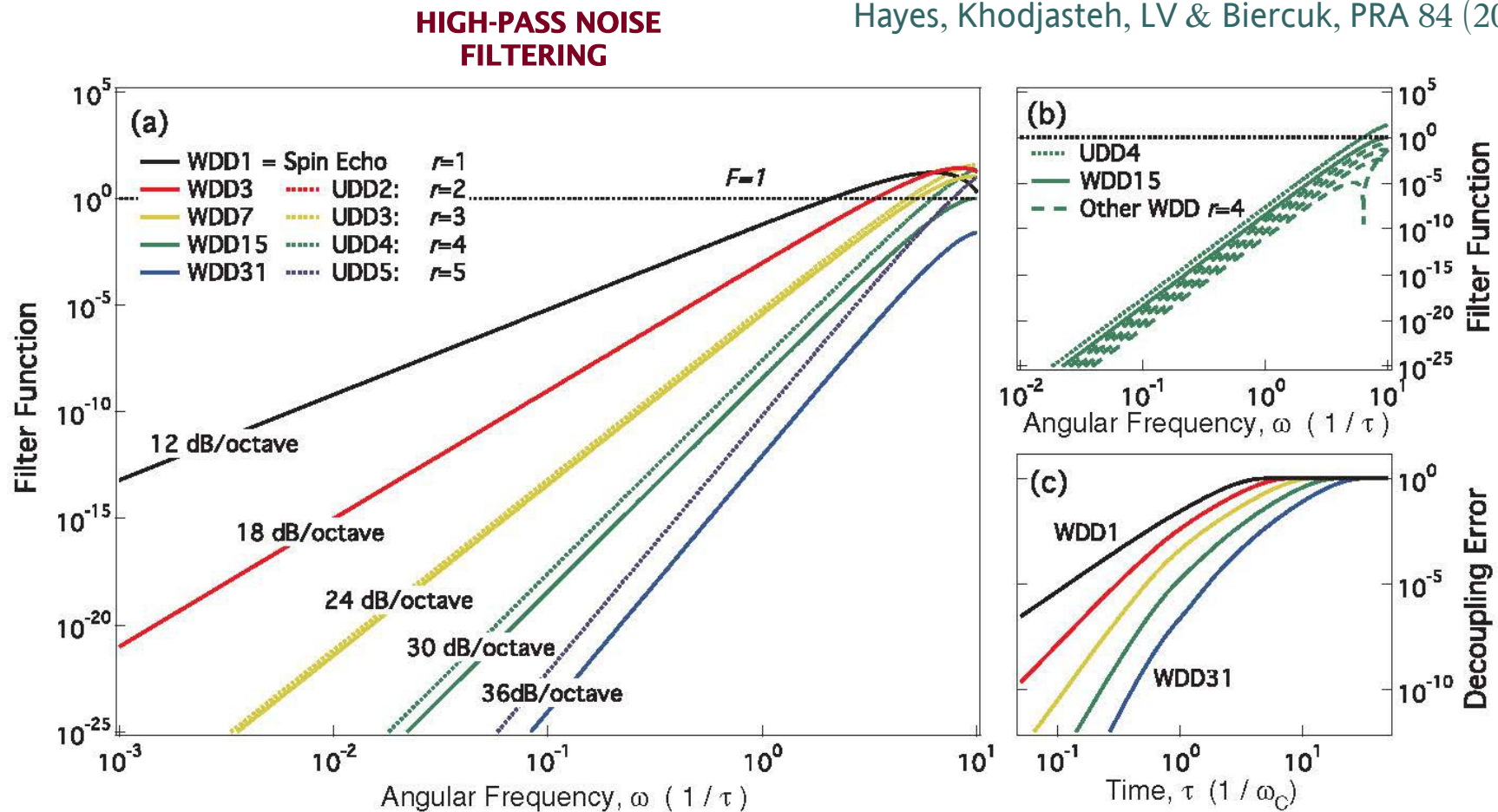
→ The larger the *order of error suppression*  $\delta$ , the higher the degree of noise cancellation:

$$\rho(t) \approx \rho(0) + \mathcal{O}(t^{\delta+1})$$



# Filter transfer function approach: Advantages...

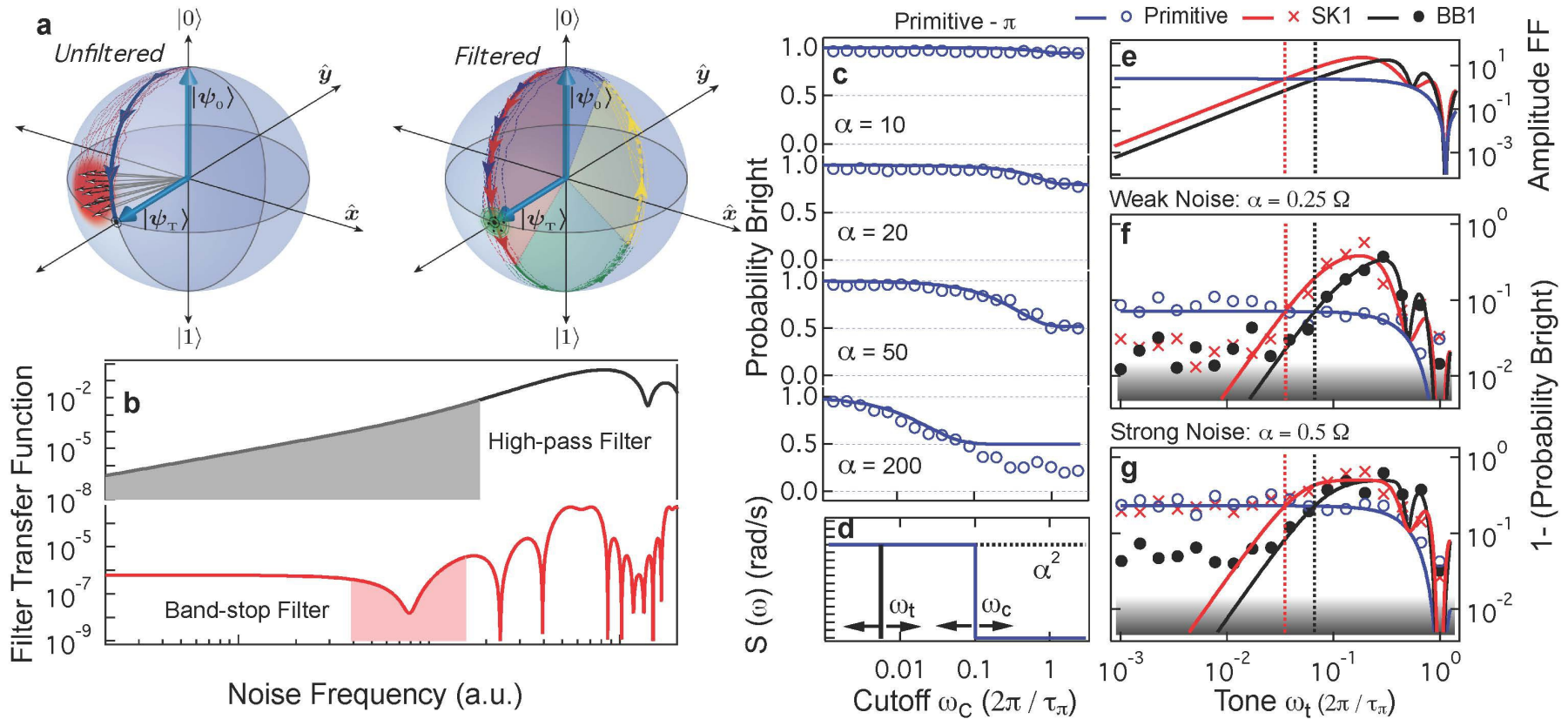
Hayes, Khodjasteh, LV & Biercuk, PRA 84 (2011).



- Direct contact with signal processing, [classical and quantum] *control engineering*...
- Simple *analytical* evaluation of control performance, compared to numerical simulation...
- Natural starting point for analysis and synthesis of control protocols *tailored to specific spectral features of generic time-dependent noise*...

# Filter transfer function approach: Validation...

Soare *et al*, Nature Phys. (Oct 2014).



- Control objective: noise-suppressed single-qubit  $\pi$  rotations under [non-Markovian] amplitude control noise  $\Rightarrow$  **Generalized FF formalism.** Green *et al*, PRL 2012, NJP 2013.
- Control protocols: [NMR] composite-pulse sequences.
- **Quantitative agreement with analytical FF predictions observed in the weak-noise limit.**

# Filter transfer function approach: Assessment...

---

- Major limitation of current generalized FF (GFF) formalism:  
High-order GFFs are given in terms of an infinite recursive hierarchy – awkward!
  - Explicit calculations to date  $\Rightarrow$  Single-qubit controlled dynamics under *classical* noise: *lowest-order* fidelity estimates, *Gaussian* [stationary] noise statistics...
  - ...
  - *Higher-order* terms are [already] of relevance to quantum control experiments...
  - What about *general* [quantum and/or non-Gaussian] noise models?...
  - What about *general target* [multi-qubit] systems?...

# Filter transfer function approach: Next steps...

- Major limitation of current generalized FF (GFF) formalism:  
High-order GFFs are given in terms of an infinite recursive hierarchy – awkward!
  - Explicit calculations to date  $\Rightarrow$  Single-qubit controlled dynamics under *classical* noise: *lowest-order* fidelity estimates, *Gaussian* [stationary] noise statistics...
  - ...
  - *Higher-order* terms are [already] of relevance to quantum control experiments...
  - What about *general* [quantum and/or non-Gaussian] noise models?...
  - What about *general target* [multi-qubit] systems?...
- Assuming that a *general* frequency-domain description is viable, to what extent will it be *equivalent* to the time-domain description...
  - How to rigorously characterize the *filtering capabilities* of a control protocol?...

## Challenge:

To build a general theory for *open-loop noise filtering* in non-Markovian quantum systems.

# Control-theoretic setting: System and noise



- Target system  $S$  (finite-dim) coupled to *quantum or classical* environment [bath]  $B$ :

$$H(t) = H_S \otimes \mathbb{I}_B + H_{SB}(t)$$

with respect to interaction picture defined by  $\mathbb{I}_S \otimes H_B$ .

→ Classical noise formally recovered for  $H(t) \equiv H_S(t)$  [stochastic time-dependence]

- Environment  $B$  is uncontrollable  $\Rightarrow$  Controller acts directly on  $S$  *alone*:

$$H_{\text{tot}} = H(t) + H_{\text{ctrl}}(t) \otimes \mathbb{I}_B \equiv \underbrace{H_0(t)}_{\text{error-free}} + \underbrace{H_e(t)}_{\text{unwanted}}$$

$$H_e(t) = H_{SB}(t) + H_{\text{ctrl},e}(t) \otimes \mathbb{I}_B + H_{S,e} \otimes \mathbb{I}_B$$

→ Evolution under ideal Hamiltonian over time  $T$  yields the desired unitary gate  $Q$  on  $S$  (e.g.,  $Q = \mathbb{I}_S$  for DD).

# Control-theoretic setting: Isolating the noise



- Total [joint] propagator may be *exactly* expressed in terms of 'error propagator':

$$U(T) = U_0(T)\tilde{U}_e(T) \equiv Q\tilde{U}_e(T) \quad \tilde{U}_e(t) = \mathcal{T}\exp\left\{-i\int_0^t \tilde{H}_e(s)ds\right\}$$

→ Choose an Hermitian operator basis on  $S$ ,  $\{\mathbb{I}_S, O_v\}$ ,  $\text{Tr}[O_v] = 0 \Rightarrow$

$$\tilde{H}_e(t) = U_0(t)^\dagger H_e(t) U_0(t) \equiv \sum_{u,v} y_{uv}(t) O_v \otimes B_u(t)$$

target-dependent  
control matrix

- Error propagator may be formally computed via a **Magnus series expansion**:

$$\tilde{U}_e(T) = e^{\sum_{\alpha=1}^{\infty} \Omega_\alpha(T)} \equiv e^{-iT(H_{SB}^{\text{eff}}(T) + H_B^{\text{eff}}(T))} \equiv e^{-i\Omega_e(T)}$$

→  $\alpha$ -th order Magnus term  $\Omega_\alpha(T)$  involves  $\alpha$ -th order nested commutators of  $\tilde{H}_e(t_j)$ .

# Cancellation order in time domain

- Magnus series has traditionally been used to characterize error-suppression properties of a control protocol in the *time domain*:

$$\tilde{U}_e(T) = e^{\sum_{\alpha=1}^{\infty} \Omega_{\alpha}(T)} \equiv e^{-iT(H_{SB}^{\text{eff}}(T) + H_B^{\text{eff}}(T))} \equiv e^{-i\Omega_e(T)}$$

Strategy: [perturbatively] **minimize the sensitivity** of the controlled evolution to  $H_e(t)$ , by making  $\tilde{U}_e(T)$  *as close as possible* to a 'pure-bath' evolution [identity on  $S$ ...]

Khodjasteh, Lidar & LV, PRL 2010; Khodjasteh, Bluhm & LV, PRA 2012.

- Definition. A control protocol specified by  $y_{uv}(t)$  achieves **cancellation order (CO)  $\delta$**  if the norm of the error action operator  $\Omega_e(T)$  [up to pure-bath terms] is reduced, such that the leading-order correction mixing  $S$  and  $B$  scales as

$$\|TH_{SB}^{\text{eff}}(T)\| = \mathcal{O}(T^{\delta+1})$$

→ CO = Standard 'decoupling order' for a DD protocol (e.g., CDD, WDD, UDD...)

# Generalized filter functions-1

- GFFs may be most generally defined directly *at the level of the effective Hamiltonian*:

→ Express each  $\tilde{H}_e(t_j)$  in the  $\alpha$ -th order term wrto the chosen operator basis:

$$\Omega_\alpha(T) = \sum_{\vec{u}, \vec{v}} \int_{\{0 \leq t_\alpha \dots \leq t_2 \leq t_1 \leq T\}} f(\{y_{[\alpha]}\}) O_{v_1} \cdots O_{v_\alpha} \otimes B_{u_1}(t_1) \cdots B_{u_\alpha}(t_\alpha) d^\alpha \vec{t}$$

→ Express each bath variable in terms of corresponding *frequency-Fourier transform*:

$$B_u(t) \equiv \int_{-\infty}^{\infty} \frac{d\omega}{2\pi} e^{i\omega t} B_u(\omega)$$



# Generalized filter functions-1

- GFFs may be most generally defined directly *at the level of the effective Hamiltonian*:

→ Express each  $\tilde{H}_e(t_j)$  in the  $\alpha$ -th order term wrto the chosen operator basis:

$$\Omega_\alpha(T) = \sum_{\vec{u}, \vec{v}} \int_{\{0 \leq t_\alpha \dots \leq t_2 \leq t_1 \leq T\}} f(\{y_{[\alpha]}\}) O_{v_1} \cdots O_{v_\alpha} \otimes B_{u_1}(t_1) \cdots B_{u_\alpha}(t_\alpha) d^\alpha \vec{t}$$

→ Express each bath variable in terms of corresponding *frequency-Fourier transform*:

$$B_u(t) \equiv \int_{-\infty}^{\infty} \frac{d\omega}{2\pi} e^{i\omega t} B_u(\omega)$$

$$\Omega_\alpha(T) = -i \sum_{\vec{u}, \vec{v}} \int \frac{d^\alpha \vec{\omega}}{(2\pi)^\alpha} G_{\vec{u}\vec{v}}^{(\alpha)}(\vec{\omega}, T) O_{v_1} \cdots O_{v_\alpha} \otimes B_{u_1}(\omega_1) \cdots B_{u_\alpha}(\omega_\alpha)$$

Meaning:  $\alpha$ -th order GFF describes the *filtering effect* of the applied control on the corresponding 'operator string' in the  $\alpha$ -th order Magnus term.

# Generalized filter functions-2

- GFFs naturally appear in the *reduced* (or *ensemble-averaged*) *system dynamics*:

→ Work in a basis where  $Q$  is diagonal and assume initial  $S$ - $B$  factorization:

$$\rho_{\ell\ell'}(T) = q_{\ell}^* q_{\ell'} \sum_{m,m'} \rho_{m,m'}(0) \text{Tr}_B [\langle \ell | \tilde{U}_e(T) | m \rangle \rho_B \langle m' | \tilde{U}_e(T)^\dagger | \ell' \rangle]$$

→ By Taylor-expanding  $\tilde{U}_e(T) = \sum_{r=0}^{\infty} (-i \Omega_e(T))^r / r!$ , and using the definition of GFFs, a *common structure* may be identified in each contributing term:

$$\langle \ell | O_{[\alpha_1]} \cdots O_{[\alpha_r]} | m \rangle \langle m' | O_{[\alpha'_1]} \cdots O_{[\alpha'_{r'}]} | \ell' \rangle$$

$$\int \mathcal{D}\vec{\omega} \underbrace{G^{(\alpha_1)} \cdots G^{(\alpha_r)} G^{*(\alpha'_1)} \cdots G^{*(\alpha'_{r'})}}_{\text{filtering properties}} \underbrace{\text{Tr}_B [\rho_B B_{[\alpha_1]} \cdots B_{[\alpha_r]} B_{[\alpha'_1]} \cdots B_{[\alpha'_{r'}]}]}_{\text{noise properties}}$$

⇒ related to high-order noise power spectra

# Generalized filter functions-2

- GFFs naturally appear in the *reduced* (or *ensemble-averaged*) *system dynamics*:

→ Work in a basis where  $Q$  is diagonal and assume initial  $S$ - $B$  factorization:

$$\rho_{\ell\ell'}(T) = q_{\ell}^* q_{\ell'} \sum_{m,m'} \rho_{m,m'}(0) \text{Tr}_B [\langle \ell | \tilde{U}_e(T) | m \rangle \rho_B \langle m' | \tilde{U}_e(T)^\dagger | \ell' \rangle]$$

→ By Taylor-expanding  $\tilde{U}_e(T) = \sum_{r=0}^{\infty} (-i \Omega_e(T))^r / r!$ , and using the definition of GFFs, a *common structure* may be identified in each contributing term:

$$\langle \ell | O_{[\alpha_1]} \cdots O_{[\alpha_r]} | m \rangle \langle m' | O_{[\alpha'_1]} \cdots O_{[\alpha'_{r'}]} | \ell' \rangle$$

$$\int \mathcal{D}\vec{\omega} \underbrace{G^{(\alpha_1)} \cdots G^{(\alpha_r)} G^{*(\alpha'_1)} \cdots G^{*(\alpha'_{r'})}}_{\text{filtering properties}} \underbrace{\text{Tr}_B [\rho_B B_{[\alpha_1]} \cdots B_{[\alpha_r]} B_{[\alpha'_1]} \cdots B_{[\alpha'_{r'}]}]}_{\text{noise properties}}$$

⇒ related to high-order noise power spectra

Example:

BB DD of a single-qubit under *Gaussian*, stationary dephasing noise [again!]

$$\chi(T) = \frac{1}{2\pi} \int_{-\infty}^{\infty} G^{(1)}(\omega, T) G^{(1)}(-\omega, T) S(\omega) d\omega, \quad G^{(1)}(\omega, T) = \int_0^T dt y(t) e^{i\omega t}$$

$$F(\omega t) \equiv \omega^2 G^{(1)}(\omega, t) G^{(1)}(-\omega, t)$$

# Fundamental filter functions

- Key insight: GFFs share a *common structure*, determined by [infinite in general, but] easily computable set of 'elemental' FFs  $\Rightarrow$  *fundamental filter functions* (FFFs):

$$F_{\vec{u}\vec{v}}^{(\alpha)}(\vec{\omega}, T) \equiv (-i)^\alpha \int_{\{0 \leq t_\alpha \leq \dots \leq t_1 \leq T\}} d^\alpha \vec{t} \prod_{j=1}^{\alpha} (y_{u_j v_j}(t_j) e^{i\omega_j t_j})$$

# Fundamental filter functions

- Key insight: GFFs share a *common structure*, determined by [infinite in general, but] easily computable set of 'elemental' FFs  $\Rightarrow$  *fundamental filter functions* (FFFs):

$$F_{\vec{u}\vec{v}}^{(\alpha)}(\vec{\omega}, T) \equiv (-i)^\alpha \int_{\{0 \leq t_\alpha \leq \dots \leq t_1 \leq T\}} d^\alpha \vec{t} \prod_{j=1}^{\alpha} (y_{u_j v_j}(t_j) e^{i\omega_j t_j})$$

- Theorem: Arbitrary GFFs of order  $\alpha$ ,  $\alpha = 1, \dots, \infty$ , may be exactly represented as

$$\begin{aligned} -iG_{\vec{u}\vec{v}}^{(\alpha)}(\vec{\omega}, T) &= F_{\vec{u}\vec{v}}^{(\alpha)}(\vec{\omega}, T) \\ &- \sum_{j=2}^{\alpha} \frac{(-1)^j}{j} \sum_{\sum_{r=1}^j \alpha_r = \alpha} \prod_{k=1}^j F_{\vec{u}_{[s_{k-1}, s_k]} \vec{v}_{[s_{k-1}, s_k]}}^{(\alpha_k)}(\vec{\omega}_{[s_{k-1}, s_k]}, T) \end{aligned}$$

$\rightarrow$  Proof follows from exact relationship between Magnus and Dyson series expansion.

Key point: Arbitrary high-order GFFs are *explicitly, non-recursively* computable as combinations of FFFs of same and lower order.

# Filtering order in frequency domain

Question: To what extent do FFFs characterize filtering properties of a protocol?

- Complete information about filtering behavior is encoded in principle in the set of all 'relevant' GFFs –  $O_{v_1} \cdots O_{v_\alpha} \neq \mathbb{I}_S$  in *at least one* factor [no pure-bath evolution].

→ For each GFF [FFF], define generalized [fundamental] CO and filtering order (FO) as

$$G_{\vec{u}\vec{v}}^{(\alpha)}(\vec{\omega}, T) \sim \mathcal{O}(m^{\Phi_{\vec{u}\vec{v}}^{(\alpha)}}(\vec{\omega} - \vec{\omega}_0) T^{\Delta_{\vec{u}\vec{v}}^{(\alpha)} + 1})$$

$$F_{\vec{u}\vec{v}}^{(\alpha)}(\vec{\omega}, T) \sim \mathcal{O}(p^{\phi_{\vec{u}\vec{v}}^{(\alpha)}}(\vec{\omega} - \vec{\omega}_0) T^{\delta_{\vec{u}\vec{v}}^{(\alpha)} + 1})$$

# Filtering order in frequency domain

Question: To what extent do FFFs characterize filtering properties of a protocol?

- Complete information about filtering behavior is encoded in principle in the set of all '**relevant**' GFFs –  $O_{v_1} \cdots O_{v_\alpha} \neq \mathbb{I}_S$  in *at least one* factor [no pure-bath evolution].

→ For each GFF [FFF], define generalized [fundamental] CO and filtering order (FO) as

$$G_{\vec{u}\vec{v}}^{(\alpha)}(\vec{\omega}, T) \sim \mathcal{O}(m^{\Phi_{\vec{u}\vec{v}}^{(\alpha)}}(\vec{\omega} - \vec{\omega}_0) T^{\Delta_{\vec{u}\vec{v}}^{(\alpha)} + 1})$$
$$F_{\vec{u}\vec{v}}^{(\alpha)}(\vec{\omega}, T) \sim \mathcal{O}(p^{\phi_{\vec{u}\vec{v}}^{(\alpha)}}(\vec{\omega} - \vec{\omega}_0) T^{\delta_{\vec{u}\vec{v}}^{(\alpha)} + 1})$$

- Definition. For a control protocol specified by  $y_{uv}(t)$ , the **generalized and fundamental cancellation order  $\Delta$  and  $\delta$**  are given by the *minimum over all the relevant GFFs/FFFs*:

$$\Delta = \min_{\mathcal{R}_\alpha, \forall \alpha} \{\Delta_{\vec{u}\vec{v}}^{(\alpha)}\}, \quad \delta = \min_{\mathcal{R}_\alpha, \forall \alpha} \{\delta_{\vec{u}\vec{v}}^{(\alpha)}\}$$

The **generalized and fundamental filtering order  $\Phi$  and  $\phi$**  at level  $\kappa$  are given by the *minimum over all the relevant GFFs/FFFs up to Magnus order  $\kappa$* :

$$\Phi^{[\kappa]} = \min_{\mathcal{R}_\alpha, \alpha \leq \kappa} \{\Phi_{\vec{u}\vec{v}}^{(\alpha)}\}, \quad \phi^{[\kappa]} = \min_{\mathcal{R}_\alpha, \alpha \leq \kappa} \{\phi_{\vec{u}\vec{v}}^{(\alpha)}\}$$

# Filtering vs. cancellation order

- Theorem: The generalized and fundamental FO and CO are related in general as follows:

$$\Phi^{[\kappa]} = \phi^{[\kappa]}, \quad \kappa = 1, \dots, \infty; \quad \Delta = \delta$$
$$\phi^{[\infty]} \leq \delta$$

Key point 1: Access to FFFs suffices to *fully* characterize the CO and FO that protocol can guarantee *under minimal assumptions on the noise model*.

- Higher effective CO and FO are possible *given specific knowledge* on the noise model.
- Level- $\kappa$  FOs are *not* a priori constrained, and the inequality at  $\kappa = \infty$  *can be strict*.

Key point 2: Cancellation and filtering are in general two *inequivalent notions*.



# Case study: Dynamical decoupling

---

- Simplest setting: **Single-axis control protocols**  $\Rightarrow$

Ideal, single-qubit DD in the presence of arbitrary, *non-Gaussian* dephasing

Claim: Arbitrarily high-order filtering may be achieved for ideal single-axis DD via concatenation,  $\text{CO} = \delta = \phi^{[\infty]} = \text{FO}$  for  $\text{CDD}_\delta$ .

$\rightarrow$  This feature is *not generic* to high-order DD protocols! E.g.  $\delta$ -th order Uhrig DD:

$$\text{CO} = \delta, \quad \text{FO} = \phi^{[\infty]} \leq 1 \text{ or } 2 \text{ for } \text{UDD}_\delta, \delta \leq 8.$$

# Case study: Dynamical decoupling

- Simplest setting: **Single-axis control protocols**  $\Rightarrow$

Ideal, single-qubit DD in the presence of arbitrary, *non-Gaussian* dephasing

**Claim:** Arbitrarily high-order filtering may be achieved for ideal single-axis DD via concatenation,  $\text{CO} = \delta = \phi^{[\infty]} = \text{FO}$  for  $\text{CDD}_\delta$ .

$\rightarrow$  This feature is *not generic* to high-order DD protocols! E.g.  $\delta$ -th order Uhrig DD:

$$\text{CO} = \delta, \quad \text{FO} = \phi^{[\infty]} \leq 1 \text{ or } 2 \text{ for } \text{UDD}_\delta, \delta \leq 8.$$

$\rightarrow$  Illustrative toy models:

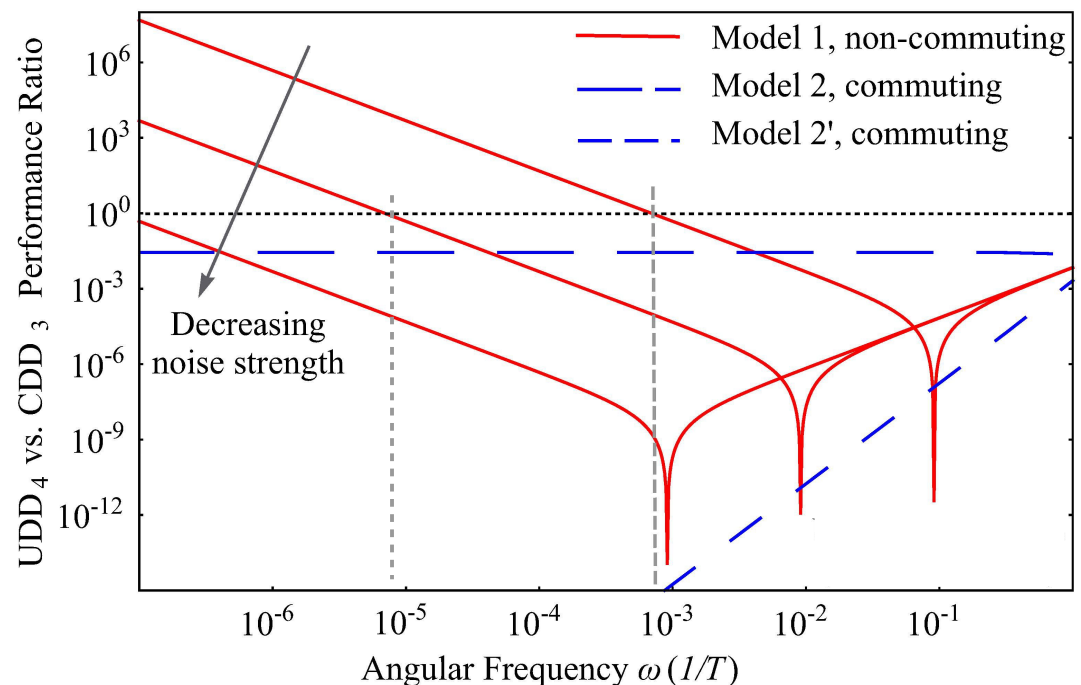
$$\tilde{H}_{e,1}(t) = gy(t)\sigma_z \otimes [B_z(t) + B_y(t)]$$

$$\tilde{H}_{e,2}(t) = gy(t)\sigma_z \otimes B_z(t)$$

$$\text{CDD}_3: \text{CO} = 3, \text{FO} = 3$$

$$\text{UDD}_4: \text{CO} = 4, \text{FO} = 2$$

Inversion of performance at low frequencies, due to high-order Magnus terms



# Further examples

---

- General case: **Multi-axis control protocols**

E.g., DD with imperfect/bounded control, DCGs, composite pulses...

Claim: A protocol which does not achieve perfect cancellation of arbitrary quasi-static noise has vanishing FO,  $\phi^{[\infty]} = 0$ .

Meaning: Arbitrarily high-order filtering is too strong a requirement – *finite- $\kappa$  filtering is relevant in practice.*

# Further examples

Soare *et al*, Nature Phys. (Oct 2014).

- General case: **Multi-axis control protocols**

E.g., DD with imperfect/bounded control, DCGs, composite pulses...

**Claim:** A protocol which does not achieve perfect cancellation of arbitrary quasi-static noise has vanishing FO,  $\phi^{[\infty]} = 0$ .

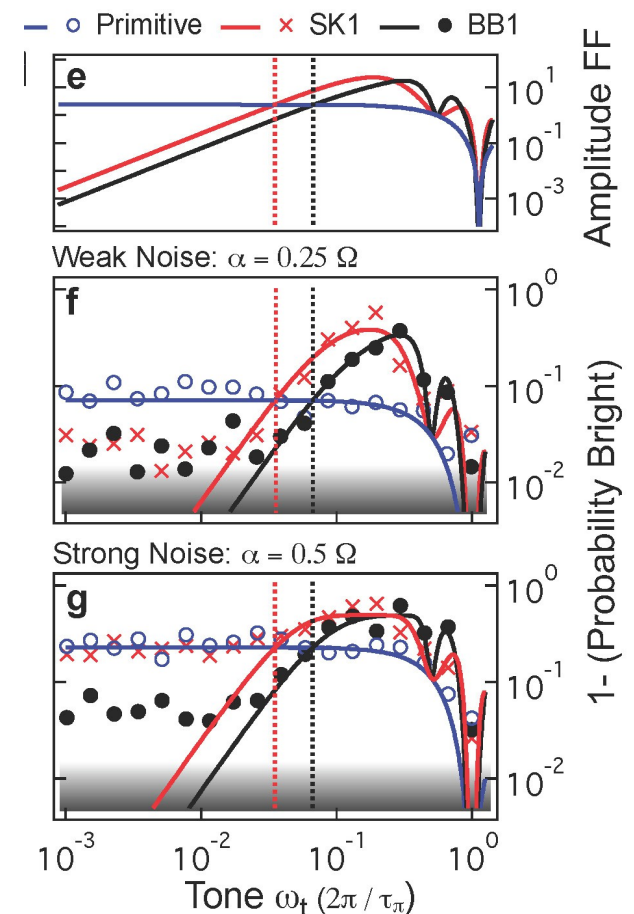
**Meaning:** Arbitrarily high-order filtering is too strong a requirement – *finite- $\kappa$  filtering is relevant in practice.*

→ Illustrative example: NMR composite-pulse sequences

SK1: CO = 1, FO = 1

BB1: CO = 2, FO = 1

→ Distinction between CO and FO is *relevant to current quantum-control experiments* and [already] informing novel approaches to control synthesis...



# Conclusion and outlook

---

- A general, *computationally tractable* approach to open-loop noise filtering in [non-Markovian] open quantum systems is possible based on identifying a set of *fundamental FFs* – out of which *arbitrary* generalized FFs may be directly assembled.
- Fundamental FFs *suffice to characterize the error-suppression capabilities in both the time and frequency domain* under minimal assumptions on the noise model.
- Order of error cancellation [a-la-Magnus] and order of filtering are in general *two inequivalent and potentially equally relevant notions for time-dependent noise*.

# Conclusion and outlook

- A general, *computationally tractable* approach to open-loop noise filtering in [non-Markovian] open quantum systems is possible based on identifying a set of *fundamental FFs* – out of which *arbitrary* generalized FFs may be directly assembled.
- Fundamental FFs *suffice to characterize the error-suppression capabilities in both the time and frequency domain* under minimal assumptions on the noise model.
- Order of error cancellation [a-la-Magnus] and order of filtering are in general *two inequivalent and potentially equally relevant notions for time-dependent noise*.
- Additional investigation is needed to appreciate the *full theoretical and experimental significance* of filtering perspective for open-loop quantum control:
  - Multi-qubit DD/long-time quantum-memory settings;  
Paz-Silva, S.-W. Lee, T. J. Green & LV, forthcoming.
  - Analytical and/or numerical synthesis of 'customized' noise filters;
  - Protocols for *non-Gaussian* noise identification/sensing;  
Paz-Silva, L. Norris & LV, forthcoming.
  - Implications for [non-Markovian] quantum fault tolerance?...
  - ⋮

