Sparse Quantum Codes from Quantum Circuits

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Joint work with D Bacon, A W Harrow, and J Shi arxiv:1411.3334

Quantum Error Correction

- Quantum error correction allows us to deal with the inevitable presence of noise in a quantum computation
- Most quantum codes are stabilizer codes.
 Ex: n=4 code that detects any single-qubit Pauli error

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4 qubits:
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Stabilizers:

 $S_X = \begin{array}{cc} X & X \\ X & X \end{array}$

 $S_Z = \begin{array}{cc} Z & Z \\ Z & Z \end{array}$

Logical qubits:

- $L_X^1 = \begin{array}{ccc} X & X \\ I & I \end{array} \qquad \begin{array}{ccc} L_X^2 = \begin{array}{ccc} X & I \\ X & I \end{array} \\ X & I \end{array}$
- $L_Z^1 = \begin{array}{ccc} Z & I \\ Z & I \end{array} \qquad \qquad L_Z^2 = \begin{array}{ccc} Z & Z \\ I & I \end{array}$

Subsystem Codes

- Solution State State
- Subsystem codes can be sparser, implying simpler syndrome measurements, higher thresholds (sometimes)

4 qubits:

$$S_{X} = G_{X}^{1} G_{X}^{2}$$

$$S_{Z} = G_{Z}^{1} G_{Z}^{2}$$
Gauge:

$$C_{X}^{1} = \begin{array}{c}X & I \\ X & I \end{array} \quad G_{X}^{2} = \begin{array}{c}I & X \\ I & X \end{array} \quad L_{X} = \begin{array}{c}X & X \\ I & I \end{array}$$

$$C_{X}^{1} = \begin{array}{c}Z & Z \\ I & I \end{array} \quad G_{X}^{2} = \begin{array}{c}I & I \\ Z & Z \end{array} \quad L_{Z} = \begin{array}{c}Z & I \\ Z & I \end{array}$$

The Structure of Subsystem Codes

- Code is defined by a set of gauge generators
- Center Z(G) of the gauge group is the stabilizer group (for a choice of signs)
- Logical operators commute with G and permute the code space (normalizer N(G))
- Bare logical operators act trivially on the gauge qubits;
 dressed ones act nontrivially



Poulin 2005

Sparse Codes

- A code family is called [n,k,d] if it encodes k logical qubits into n physical qubits and can detect any Pauli error of weight < d.</p>
- A code with a given set of gauge generators is called s-sparse if:
 - \ll every gauge generator has weight \leq s

 - % ex: row- and column-sparse parity-check matrix
- * A code is called just **sparse** if s = O(1), independent of *n*

The Importance of Sparse Codes

- Only at most s qubits need to be measured at a time, instead of O(n)
- > higher thresholds, parallelized architectures, simpler decoding algorithms, FTQC with low overhead
- Ex: topological codes; LDPC codes
 - * toric code, color codes, hypergraph product codes, ...
- A major challenge is to find sparse quantum codes that perform well, e.g. with k, d = O(n) and fast decoders



Gallager 1962, MacKay & Neal 1995, Kitaev 2003, Dennis *et al.* 2001, Raussendorf & Harrington 2007, Tillich & Zémor 2009, Kovalev & Pryadko 2013, Bravyi & Hastings 2013, Gottesman 2013, ...

Main Result

Theorem 1. Given any $[n_0, k_0, d_0]$ quantum stabilizer code with stabilizer generators of weight $w_1, \ldots, w_{n_0-k_0}$, there is an associated [n, k, d] quantum subsystem code whose gauge generators have weight O(1) and where $k = k_0$, $d = d_0$, and $n = O(n_0 + \sum_i w_i)$. This mapping is constructive given the stabilizer generators of the base code.

A systematic way to convert *any* stabilizer code into a **sparse** subsystem code with the **same** *k* and *d* parameters

The **price** is an increase in the number of physical qubits equal to the sum of the original generator weights

The proof is hard, but the construction itself is quite simple

- Begin with a stabilizer code of your choice
- Write a quantum circuit for measuring the stabilizers of this code.



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- Begin with a stabilizer code of your choice
- Write a quantum circuit for measuring the stabilizers of this code.
- Turn the circuit elements into input/output qubits
- Add gauge generators via Pauli circuit identities
- This defines the code



Properties of this Construction

Circuits as linear operators preserving the code space



$$V = |00\rangle\!\langle 00| + |11\rangle\!\langle 11|$$

 $C = \operatorname{span}(\{|00\rangle, |11\rangle\})$ V is a good circuit

> General condition: V is good iff $V^{\dagger}V = \Pi_{\mathcal{C}}$

Properties of this Construction

- Circuits as linear operators preserving the code space



Apply gauge operators...

Properties of this Construction

- Circuits as linear operators preserving the code space
- Squeegee lemma: using gauge operations, we can localize errors to the initial data qubits



Stabilizer and Logical Operators

- Spackling: like squeegee, but you leave a residue
- Spackling of logical
 operators gives the new
 logical operators
- Spackling of stabilizers on the inputs and ancillas are the new stabilizers
- Everything else is gauge or detectable error
- * ...what about distance?



	Λ	Λ	Λ		\boldsymbol{Z}	\Box	\Box
$L_X =$	X	X	X	$L_Z =$	Ι	Ι	Ι
	Ι	X	Ι		Ι	Ι	Ι
	Z	Z	Z		Z	Z	Ι
S =	Z	Z	Z	$S_a =$	Z	Ι	Ι
	I	Ι	I		Z	Z	Z

*even/odd effect means that circuits wires must have odd length

Code Distance and Fault Tolerance

- For most syndrome-measurement circuits, the new code is uninteresting
- If we use a fault-tolerant circuit then we preserve the code distance
- **Solution** Fault tolerance: for every error pattern *E*, either $V_E = 0$ or there exists *E*' on inputs s.t. *V E*'=*V*_{*E*} and |E'| ≤ |E|
- Strange constraints:
 - Circuit must be Clifford (so no majority vote)
 - No classical feedback or post-processing allowed
 - However, we only need to detect errors

Fault-Tolerant Gadgets

- Se modified Shor/ Divincenzo cat states
- Build a cat, and postselect ...not fault tolerant
- Redeem this idea by coupling to expanders
- Constant-degree expanders exist with sufficient edge expansion to make this fault tolerant



Theorem 1. Given any $[n_0, k_0, d_0]$ quantum stabilizer code with stabilizer generators of weight $w_1, \ldots, w_{n_0-k_0}$, there is an associated [n, k, d] quantum subsystem code whose gauge generators have weight O(1) and where $k = k_0$, $d = d_0$, and $n = O(n_0 + \sum_i w_i)$. This mapping is constructive given the stabilizer generators of the base code.

- Created sparse subsystem codes with the same k and d parameters as the base code
- Sed fault-tolerant circuits in a new way, via expanders
- Extra ancillas are required according to the circuit size

Almost "Good" Sparse Subsystem Codes

- Start with an $[n_0, 1, d_0]$ random stabilizer code (so that $d_0 = O(n_0)$ with high probability)
- * Concatenate this *m* times to get an $[n_0^m, 1, d_0^m]$ code
- Sum over the stabilizer weights gives $n = n_0^{O(m)}$
- **Solution** & Apply Theorem 1 with $m = O(\sqrt{\log n})$

Sparse subsystem codes exist with $d = O(n^{1-\varepsilon})$ and $\varepsilon = O(1/\sqrt{\log n})$.

Best previous distance for sparse codes was $d = O(\sqrt{n \log n})$ by Freedman, Meyer, Luo 2002

*Thank you Sergei Bravyi!

Local Subsystem Codes Without Strings

- Take the circuit construction from the previous result
- SWAP gates and wires, spread the circuit over the vertices of a cubic lattice in D dimensions
- * Let $n=L^D$ be the total number of qubits

Local subsystem codes exist with $d = O(L^{D-1-\varepsilon})$ and $\varepsilon = O(1/\sqrt{\log n})$.

Compared to Known Bounds

- * Our code: $d = \Omega(L^{D-1-\varepsilon})$
- ❀ Best known local stabilizer codes: d=O(L^{D/2})
- Local commuting projector codes
 $kd^{2/(D-1)} \le O(n)$
 - * Our codes: $kd^{2/(D-1)} = \Omega(n)$ (use the hypergraph product codes and Thm 1)

Bravyi & Terhal 2009; Bravyi, Poulin, Terhal 2010; Tillich & Zémor 2009

 $\varepsilon = O(1/\sqrt{100})$

Local Subsystem Codes Without Strings

- Specialize to D=3
- * Sparse subsystem code on a lattice with $[L^3, O(1), L^{2-\varepsilon}]$
- No strings, either for bare or dressed logical operators
 cf. Bombin's gauge color codes
- …on the other hand it's a subsystem code
- How does this compare to other candidate selfcorrecting quantum memories?



Comparing Candidate Self-Correcting Memories

Code	Self-correcting?	Comments
3D Bacon-Shor (Bacon 2005)	no	No threshold, so no self- correction (Pastawski <i>et al</i> . 2009)
Welded Code (Michnicki 2014)	no	See Brown <i>et al</i> . 2014 review article for discussion
Cubic Code (Haah 2011)	marginal	poly(L) memory lifetime for L< e ^{β/3} (Bravyi & Haah 2013)
Embedded Fractal Product Codes (Brell 2014)	maybe	very large ground-state degeneracy?
Gauge Color Codes (Bombin 2013)	???	Does have a threshold, also has string-like dressed operators
This talk (BFHS 2014)	???	No strings, concatenated codes have a threshold

Not depicted: Codes with long-range couplings (e.g. several works by the Loss group) or Hamma *et al.* 2009 See the talk by Olivier Landon-Cardinal on Friday for more discussion of these types of codes.

Challenges with Gauge Hamiltonians

- Gauge Hamiltonians are sometimes gapped: (Kitaev 2005; Brell et al. 2011; Bravyi et al. 2013)
- …but sometimes not: (Bacon 2005; Dorier, Becca, & Mila 2005)
- The simplest example of our code (a wire) reduces to Kitaev's quantum wire, which is gapped as long as the couplings aren't equal in magnitude
- Our codes are a vast generalization of Kitaev's wire to arbitrary circuits!
- This undoubtedly has a rich phase diagram... might there be a gapped self-correcting phase, or something more?

Kitaev 2001; Lieb, Schultz, & Mattis 1961

Conclusion & Open Questions

- Showed a generic way to turn stabilizer codes into sparse subsystem codes
- New connection between quantum error correction & fault-tolerant quantum circuits
- What are the limits for sparse stabilizer codes?
- Self-correcting memory from the gauge Hamiltonian?
- Setting the set of the set of
- Improve the rate? (Bravyi & Hastings 2013)
- Section 2 Sec
- See arxiv:1411.3334 for more details!